

Polyakov Loop Susceptibility and Correlators in the Chiral Limit

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Lattice 2019, 17 June 2019



Polyakov loop correlators

The untraced Polyakov loop

$$L_{\vec{x}} = \prod_{\tau} U_4(\vec{x}, \tau)$$

is related to the color-averaged free energy of a quark-antiquark pair¹

$$F_{q\bar{q}}(r, T) = -T \log \left\langle \frac{1}{9} \text{tr } L_{\vec{x}} \text{tr } L_{\vec{y}}^\dagger \right\rangle \quad r = |\vec{x} - \vec{y}|.$$

Of interest to us will be color-singlet free energy²

$$F_1(r, T) = -T \log \left\langle \frac{1}{3} \text{tr } L_{\vec{x}} L_{\vec{y}}^\dagger \right\rangle,$$

which is not a gauge invariant quantity.

¹L. D. McLerran and B. Svetitsky, Phys. Rev. D, 24.2, 450–460 (1981).

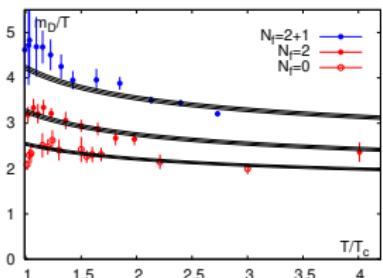
²S. Nadkarni, Phys. Rev. D, 34.12, 3904–3911 (1986).

Debye mass m_D

- $r_D = 1/m_D$ characterizes distance at which in-medium modifications of quark-antiquark interaction dominate (**color screening**).
- Can extract m_D from large r behavior of F_1 :

$$F_1(r, T) \sim \frac{\alpha(T)}{r} e^{-r m_D(T)} + C$$

- m_D dependence on T and N_f can be seen, e.g. in lattice simulations³.



QUESTION: How does m_D depend on m_ℓ ?

³O. Kaczmarek, PoS(CPOD07), 043 (2008).

The Polyakov loop,

$$P = \frac{1}{N_s^3} \sum_{\vec{x}} \frac{1}{N_c} \text{tr } L_{\vec{x}},$$

m_D , and deconfinement are all related.

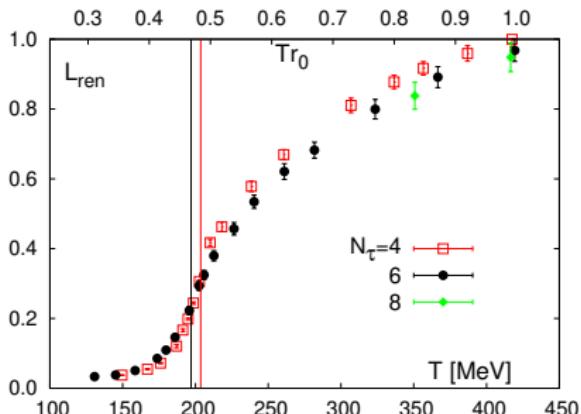
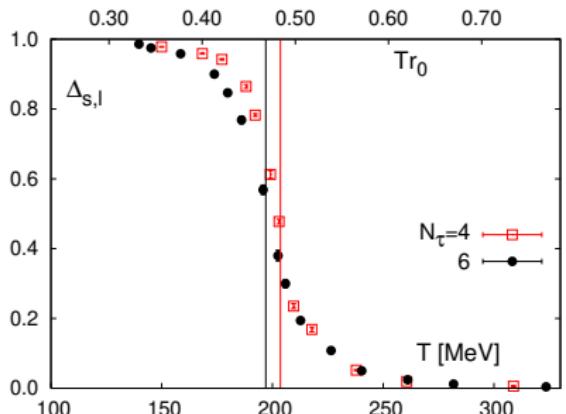
- For quenched QCD

$$\chi = N_s^3 \left(\langle |P|^2 \rangle - \langle |P| \rangle^2 \right)$$

peaks near T_c , where \mathbb{Z}_3 is spontaneously broken.

- Corresponds to an inflection point in P .
- Finite quark mass breaks \mathbb{Z}_3 explicitly.
- Nevertheless at large quark mass some remnant seemed to remain.
- Tempting to associate χ peak with hadron melting in dynamical QCD.

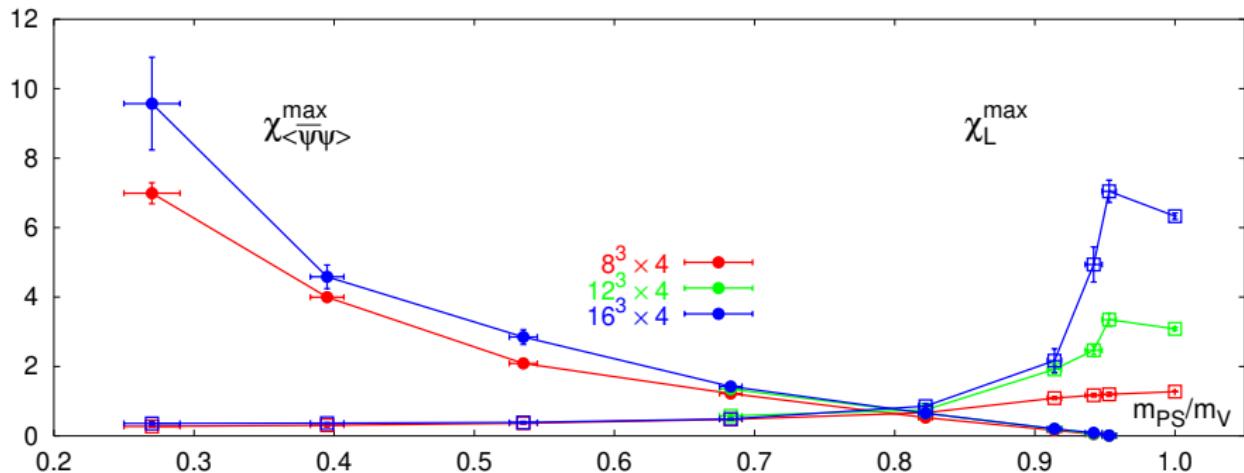
Deconfinement criteria



- Past studies have shown order parameter inflection points to appear at similar temperatures⁴ however...

⁴M. Cheng et al., Phys. Rev. D, 77.1, 014511 (2008).

Deconfinement criteria



- ...the heights of susceptibility maxima have been known for some time to depend strongly on the quark mass⁵.

QUESTION: Does it make sense to associate Polyakov susceptibility with hadron melting, especially as $m_\ell \rightarrow 0$?

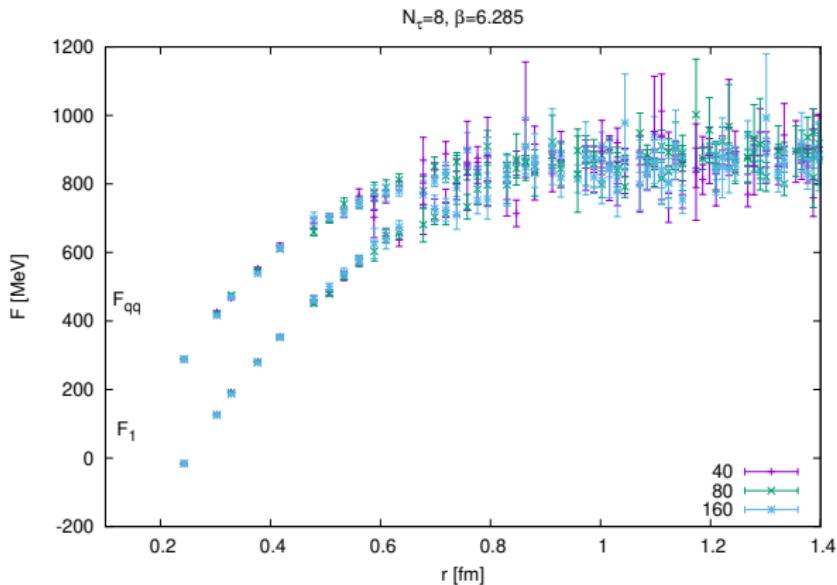
⁵F. Karsch, Lectures on Quark Matter, 583, 209–249 (2002).

- $N_f = 2 + 1$ with HISQ action
- $N_\tau = 8$ and 12
- $N_s/N_\tau \geq 3$
- m_s fixed to its physical value
- m_s/m_ℓ varies from 27 to 160
- T in the vicinity of chiral transition temperature
- Set scale with r_1 ⁶
- F_{qq} and F_1 measurements in Coulomb gauge
- Renormalize by matching F_1 to zero temperature potential⁷
- Roughly 3000 to 20000 depending on the parameters

⁶C. Bernard, PoS(Lattice 2010), 074 (2011).

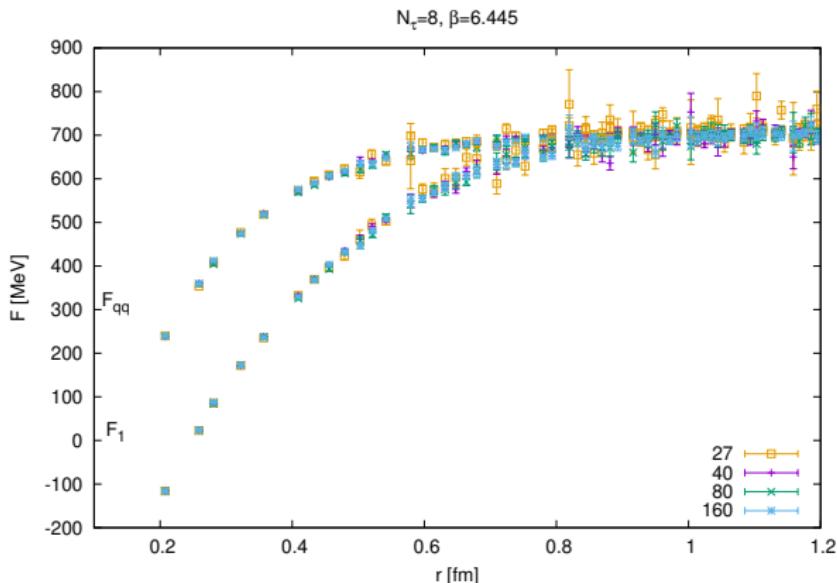
⁷O Kaczmarek et al., Phys. Lett. B, 543.1-2, 41–47 (2002).

Preliminary results: Free energies



$$T \approx 141 \text{ MeV}$$

Preliminary results: Free energies

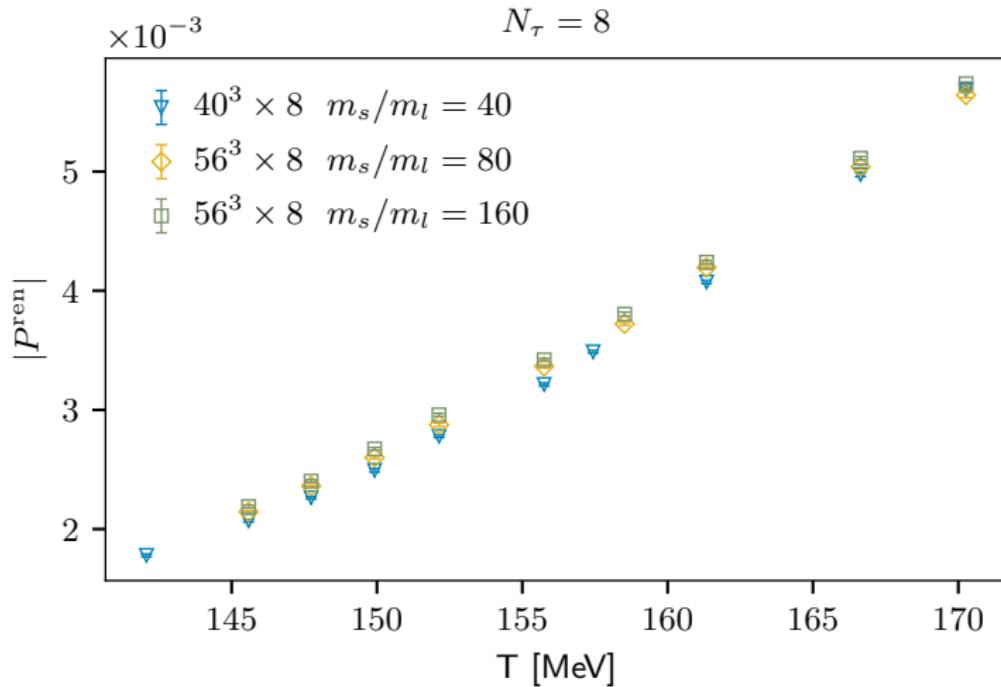


$$T \approx 166 \text{ MeV}$$

CONCLUSION: No dependence of m_D on m_ℓ noticeable.

FUTURE: Precise determination of m_D . (Gradient flow?)

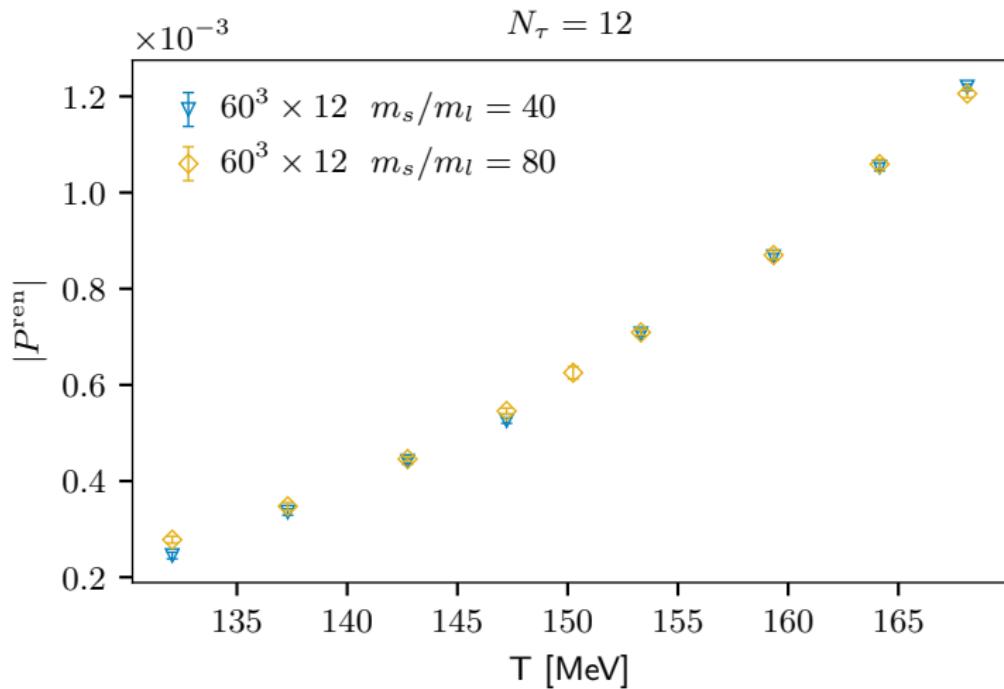
Preliminary results: Polyakov loop



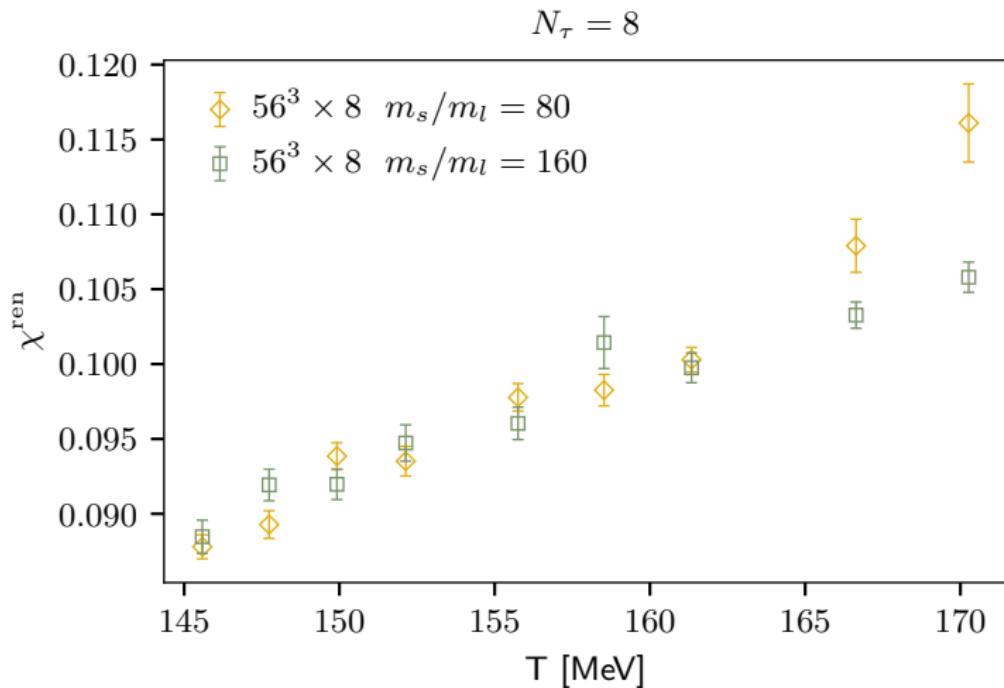
$$T_{pc}^{40} \approx 158 \text{ MeV}, \quad T_{pc}^{80} \approx 154 \text{ MeV}, \quad T_{pc}^{160} \approx 151 \text{ MeV}$$

No inflection point

Preliminary results: Polyakov loop



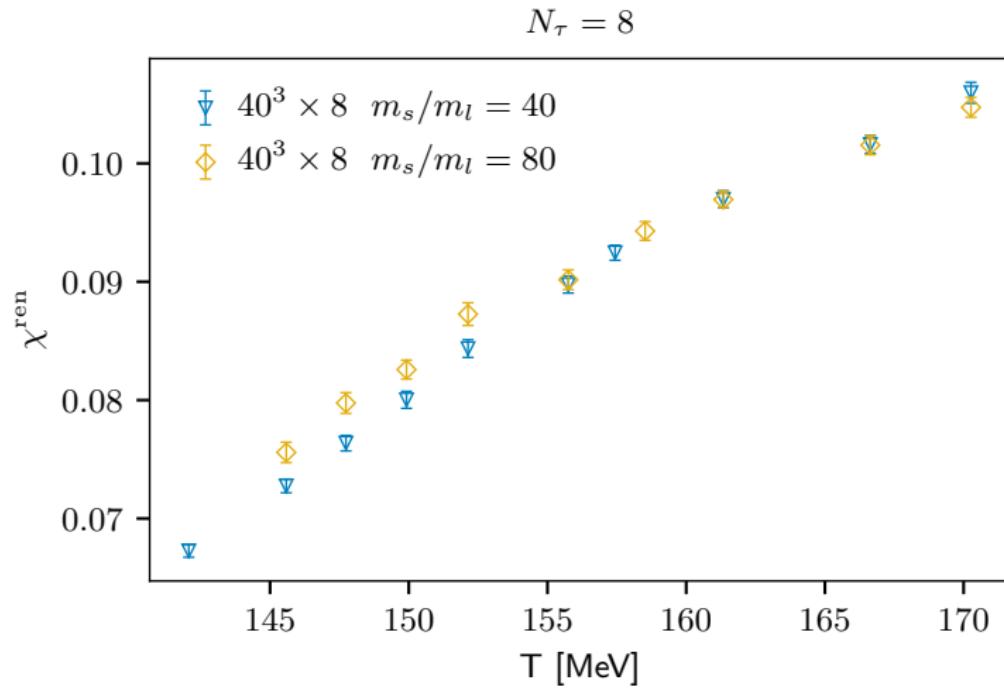
Preliminary results: Polyakov susceptibility



$$T_{pc}^{80} \approx 154 \text{ MeV}, \quad T_{pc}^{160} \approx 151 \text{ MeV}$$

No peak

Preliminary results: Polyakov susceptibility



$$T_{pc}^{40} \approx 158 \text{ MeV}, \quad T_{pc}^{80} \approx 154 \text{ MeV}$$

CONCLUSION: No drastic change in P^{ren} near chiral T_{pc} .

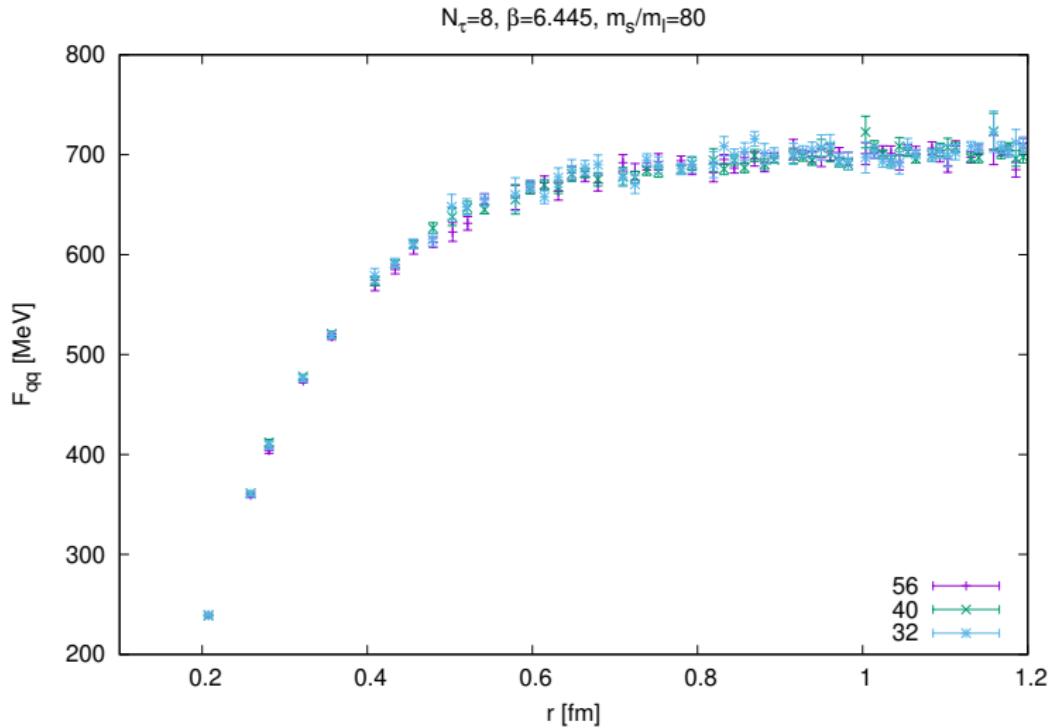
- Does m_D change with m_ℓ ?
 - Preliminary results suggest no dependence within our statistics.
 - Working toward numerical determination of m_D .
 - May use gradient flow to smooth UV fluctuations.
- χ as a probe for hadron melting?
 - Does not coincide with T_{pc} from chiral susceptibility.
 - Results at other parameter combinations still forthcoming...
 - ...in particular points at $m_s/m_l = 27$ and 20 .

Thank you!

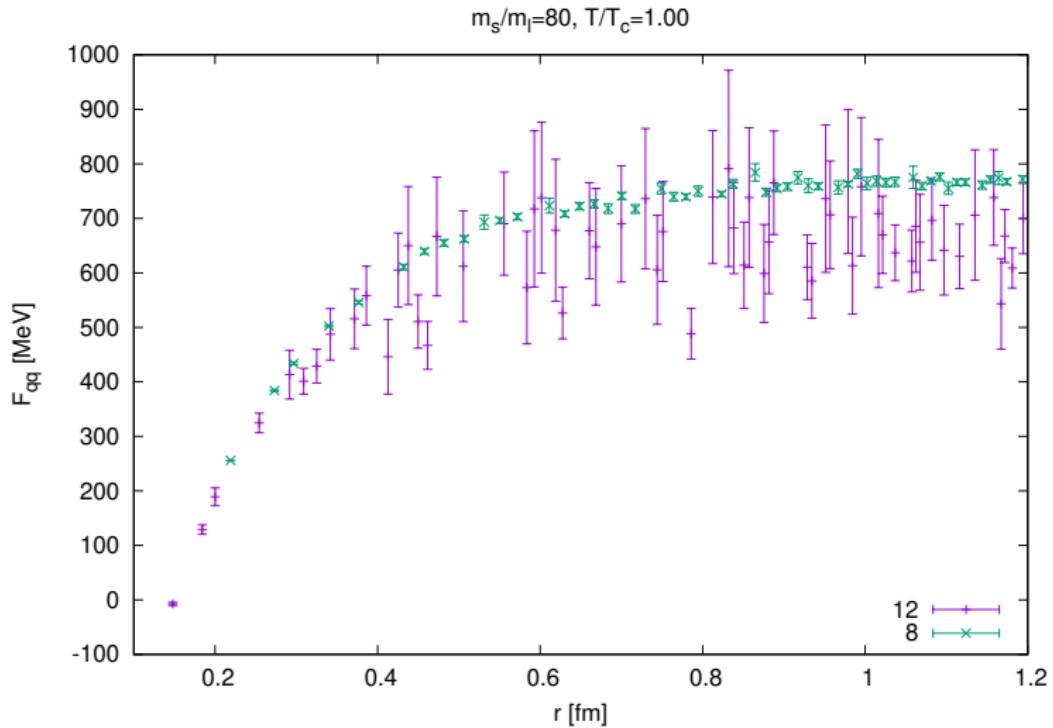
Additional details: current statistics

| $N_s^3 \times N_\tau$ | m_s/m_ℓ | approx. $N_{\text{conf}}/\text{run}$ |
|-----------------------|--------------|--------------------------------------|
| $24^3 \times 8$ | 40 | 20 000 |
| $32^3 \times 8$ | 27 | 6 000 |
| | 80 | 10 000 |
| $40^3 \times 8$ | 40 | 10 000 |
| | 80 | 6 000 |
| $42^3 \times 12$ | 40 | 10 000 |
| $48^3 \times 12$ | 80 | 6 000 |
| $56^3 \times 8$ | 160 | 3 000 |
| | 80 | 3 000 |
| $60^3 \times 12$ | 40 | 6 000 |
| | 80 | 3 000 |

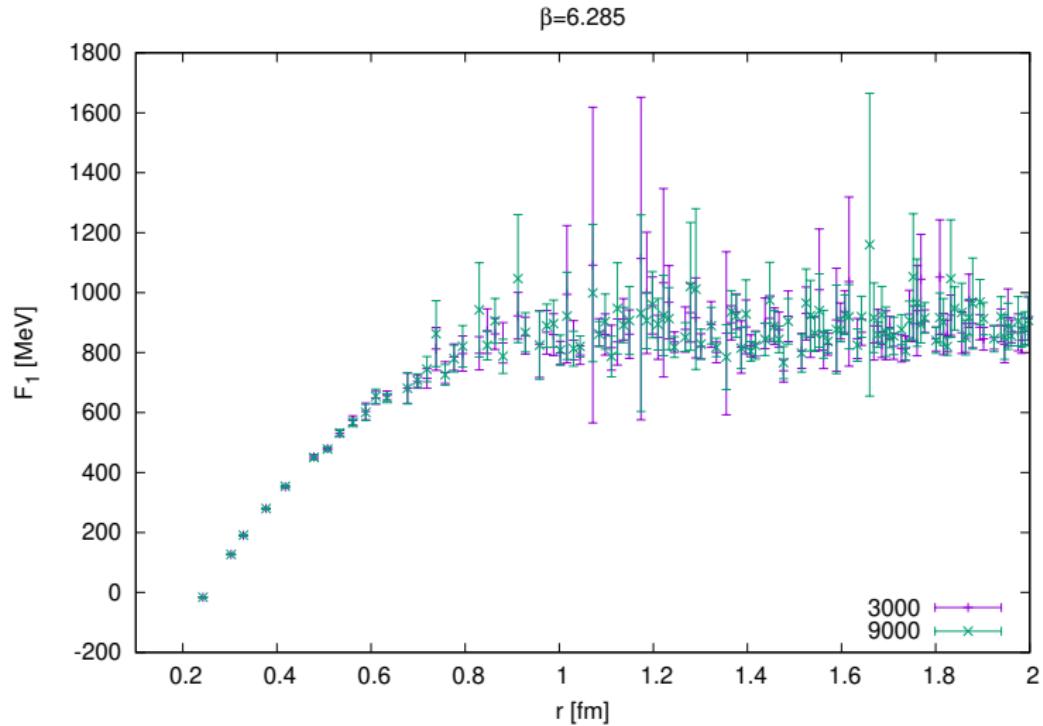
Additional details: N_s dependence of F_{qq}



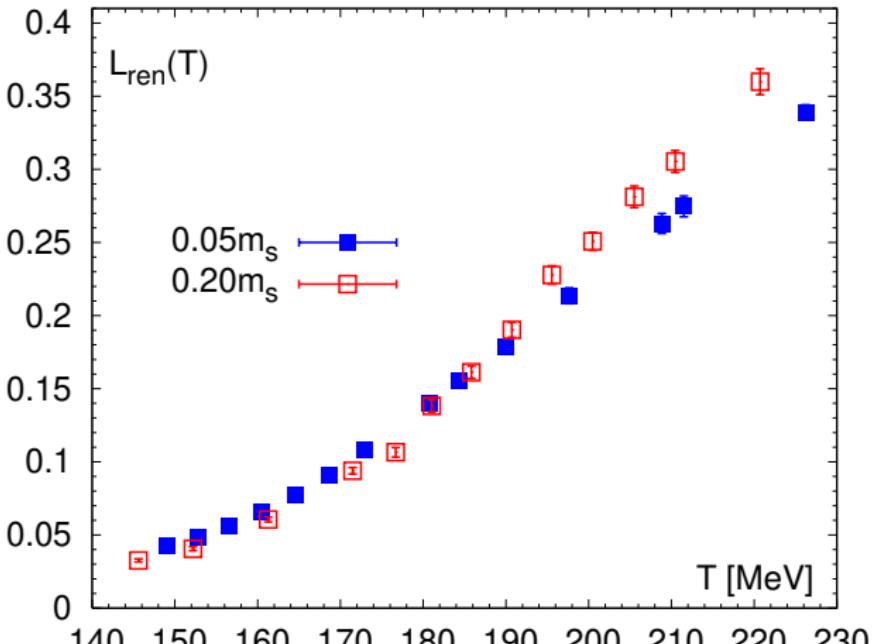
Additional details: a dependence of F_{qq}



Additional details: θ_{gf} dependence of F_1



Additional details: P inflection points



Taken from a HISQ study⁸.

⁸A. Bazavov et al., Phys. Rev. D, 85.5 (2012).